

# Consistent-Mode Indicator for the Eigensystem Realization Algorithm

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**A new method is described for assessing the consistency of structural modal parameters identified with the eigensystem realization algorithm. Identification results show varying consistency in practice due to many sources including high modal density, nonlinearity, and inadequate excitation. Consistency is considered to be a reliable indicator of accuracy. The new method is the culmination of many years of experience in developing a practical implementation of the eigensystem realization algorithm. The effectiveness of the method is illustrated using data from NASA Langley's controls-structures-interaction evolutionary model.**

## Introduction

THE dynamic behavior of most aerospace structures is adequately described using modal parameters (natural frequencies, mode shapes, damping factors, and modal masses). The objective of structural modal identification is to obtain a valid modal representation over a specified frequency range for all spatial degrees of freedom. This objective is considerably different than identifying an input-output map only at particular degrees of freedom where control actuators and sensors are located.<sup>1</sup> A full spatial modal representation permits several tasks to be performed that cannot be performed using an input-output map derived for control purposes. These tasks include validation of structural modeling procedures and assumptions, prediction of system dynamics using modal parameters of individual components, investigation of more effective actuator or sensor locations for control purposes, and improved characterization of disturbances occurring at unexpected locations on a spacecraft during operation.

It is relatively straightforward to estimate structural modal parameters experimentally using a variety of available approaches.<sup>2,3</sup> However, it is generally much more difficult to establish reliable confidence values for each result. Confidence criteria based on noise characteristics are available<sup>4</sup> but are of limited usefulness in practical applications. In modal-survey tests, identification difficulties arise primarily from high modal density, nonlinearity, weakly excited modes, local modes, nonstationarities, rattling, etc., not from instrumentation noise. The simultaneous effects of these conditions are in general impossible to include explicitly in confidence calculations.

The consistent-mode indicator (CMI) described in this paper provides a reliable, relative measure of accuracy for structural modal parameters identified with the eigensystem realization algorithm (ERA).<sup>5-7</sup> A single value ranging from 0–100% is obtained for each identified mode. Furthermore, the results can be decomposed into constituent components associated

with each input (initial condition) and output (response measurement) or input-output pair. Both temporal and spatial consistency calculations are included in the formulation. Modes with CMI values greater than approximately 80% are identified with high confidence. Modes with values ranging from 80–1% display moderate to large uncertainty. Fictitious “computational modes” have CMI values of zero.

The first part of this paper contains a brief summary of ERA followed by a complete description of CMI. The second part illustrates the concepts using recent laboratory data from NASA Langley's controls-structures-interaction (CSI) evolutionary model (CEM).<sup>8</sup> The CEM is a large flexible research structure being used to experimentally assess the level of confidence with which CSI technology can be applied to future spacecraft.

## Eigensystem Realization Algorithm

A finite-dimensional, linear, time-invariant dynamic system can be represented by the state-variable equations

$$\begin{aligned}\dot{x}(t) &= A_c x(t) + B u(t) \\ y(t) &= C x(t)\end{aligned}\quad (1)$$

where  $x$  is an  $n$ -dimensional state vector,  $u$  is a  $p$ -dimensional excitation vector, and  $y$  is a  $q$ -dimensional response vector. A special solution to these equations is the impulse response function

$$Y(t) = C e^{A_c t} B \quad (2)$$

for  $t \geq 0$ . The  $i$ th column of  $Y(t)$  contains the free response of the system, with  $x(0) = 0$ , to a unit-impulse excitation applied at the  $i$ th input location at  $t = 0$ .

With sampled data, this solution can be expressed as

$$Y(k) = C A^k B \quad (3)$$

for  $k \geq 0$ .  $A = e^{A_c \Delta t}$  is the state-transition matrix and  $\Delta t$  is the data sampling interval.

The problem of system realization is as follows: Given a sequence of experimentally measured matrices  $Y(k)$ , for  $k = 0, 1, 2, \dots$ , construct a triplet  $[A, B, C]$  such that Eq. (3) is satisfied as closely as possible. Note that  $[A, B, C]$  is not unique since the set  $[T^{-1}AT, T^{-1}B, CT]$ , for any nonsingular matrix  $T$ , also satisfies the state-variable equations.

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The ERA solution to the system realization problem uses singular value decomposition of the generalized Hankel matrix

$$H(0) = U \Sigma V^T \quad (4)$$

Matrix  $H(0)$  consists of time-shifted  $Y(k)$  submatrices. In practice, every block row and block column of  $H(0)$  normally contains data shifted in time by one sample from data in the previous block row or block column. The only exceptions are the final block row and column, which are shifted by a larger number of time samples (by default, 10 samples) to calculate the extended modal amplitude coherence (EMAC) discussed later in the paper. Also, in large modal surveys involving hundreds of response measurements, many of the rows of  $H(0)$  below row  $q$  ( $q$  = number of measurements) may generally be deleted without loss of accuracy due to the large redundancy of information.

Retaining the  $n$  largest singular values and corresponding vectors, an  $n$ th-order realization is computed as follows:

$$\begin{aligned} A &= \Sigma_n^{1/2} U_n^T H(1) V_n \Sigma_n^{-1/2} \\ B &= \Sigma_n^{1/2} V_n^T(:, p) \\ C &= U_n(q, :) \Sigma_n^{1/2} \end{aligned} \quad (5)$$

where  $H(1)$  is a matrix of the same form as  $H(0)$  but whose data are shifted in time by one additional sample. The notations  $(:, p)$  and  $(q, :)$  in Eq. (5) indicate the first  $p$  columns and first  $q$  rows, respectively, of the corresponding matrices.

This realization is transformed to modal coordinates using the eigenvalues  $Z$  and eigenvector matrix  $\Psi$  of  $A$

$$\begin{aligned} A' &= \Psi^{-1} A \Psi = Z \quad (\text{diagonal}) \\ B' &= \Psi^{-1} B \\ C' &= C \Psi \end{aligned} \quad (6)$$

The modal damping rates  $\sigma_i$  and damped natural frequencies  $\omega_i$  are the real and imaginary parts of the eigenvalues after transformation back to the continuous domain

$$s_i = \sigma_i \pm j\omega_i = \ell_n(z_i)/\Delta t \quad (7)$$

Modal participation factors and mode shapes are the corresponding rows of  $B'$  and columns of  $C'$ , respectively.

In practice, some modal parameters obtained using this approach are inaccurate due to high modal density, nonlinearity, etc. CMI is used to assess the relative accuracy of the various results.

### Consistent-Mode Indicator

CMI is computed for mode  $i$  as the product of two other parameters, EMAC and modal phase collinearity (MPC)

$$\text{CMI}_i = \text{EMAC} \cdot \text{MPC}_i \quad (\times 100\%) \quad (8)$$

EMAC quantifies the temporal consistency of the identification results. MPC quantifies the spatial consistency of the identification results. Practical experience has shown that both conditions must be satisfied simultaneously to ensure accurate results. These two parameters are described separately in the following sections.

### Extended Modal Amplitude Coherence

EMAC is computed using the identified modal parameters. Mode shape and modal participation components for data at  $t = 0$  are compared with corresponding components at  $t = T_o$  (for outputs) and  $t = T_i$  (for inputs) located in the final block row and final block column, respectively, of the modal

observability matrix  $G' = U_n \Sigma_n^{1/2} \Psi$  and modal controllability matrix  $F' = \Psi^{-1} \Sigma_n^{1/2} V_n^T$ . Data in the corresponding final block row and final block column of the Hankel matrices are shifted by 10 time samples (by default) from the previous block row and column, providing an extension of the primary data analysis window. An EMAC value is computed for each of the  $p$  inputs (initial conditions) and  $q$  outputs (response measurements), for every mode.

Let  $(\phi_{ij})_0$  be the identified mode shape component for mode  $i$  and response measurement  $j$  at  $t = 0$  and  $(\phi_{ij})_{T_o}$  be the corresponding identified component at  $t = T_o$ . The identified eigenvalue for mode  $i$  is  $s_i$ . Compute a predicted value of  $(\phi_{ij})_{T_o}$  as follows:

$$(\tilde{\phi}_{ij})_{T_o} = (\phi_{ij})_0 \cdot e^{s_i T_o} \quad (9)$$

Temporal consistency is quantified by comparing  $(\phi_{ij})_{T_o}$  and  $(\tilde{\phi}_{ij})_{T_o}$ . The actual and predicted magnitudes are compared using the ratio of the magnitudes

$$\begin{aligned} R_{ij} &= |(\phi_{ij})_{T_o}| / |(\tilde{\phi}_{ij})_{T_o}| \quad \text{for } |(\phi_{ij})_{T_o}| \leq |(\tilde{\phi}_{ij})_{T_o}| \\ &= |(\tilde{\phi}_{ij})_{T_o}| / |(\phi_{ij})_{T_o}| \quad \text{otherwise} \end{aligned} \quad (10)$$

The actual and predicted phase angles are also compared. Letting  $P_{ij} = \text{ARG}[(\phi_{ij})_{T_o} / (\tilde{\phi}_{ij})_{T_o}]$ ,  $-\pi \leq P_{ij} \leq \pi$ , a weighting factor is determined as follows:

$$\begin{aligned} W_{ij} &= 1.0 - (|P_{ij}| / (\pi/4)) \quad \text{for } |P_{ij}| \leq \pi/4 \\ &= 0.0 \quad \text{otherwise} \end{aligned} \quad (11)$$

An output EMAC for mode  $i$  and response measurement  $j$  is then computed as

$$\text{EMAC}_{ij}^o = R_{ij} \cdot W_{ij} \quad (\times 100\%) \quad (12)$$

An input EMAC for mode  $i$  and initial condition  $k$ ,  $\text{EMAC}_{ik}^I$ , is similarly computed using the identified modal participation factors.

Using these results, an EMAC value is associated with every  $j$ - $k$ th input-output pair as follows:

$$\text{EMAC}_{ijk} = \text{EMAC}_{ij}^o \cdot \text{EMAC}_{ik}^I \quad (\times 100\%) \quad (13)$$

Finally, to condense all EMAC results for mode  $i$  into a single value, a weighted average of the individual results is computed

$$\begin{aligned} \text{EMAC}_i &= \frac{\sum_{j=1}^q \sum_{k=1}^p \text{EMAC}_{ijk} \cdot |\phi_{ij}|^2 \cdot |\phi_{ik}|^2}{\sum_{j=1}^q \sum_{k=1}^p |\phi_{ij}|^2 \cdot |\phi_{ik}|^2} \\ &= \frac{\left( \sum_{j=1}^q \text{EMAC}_{ij}^o \cdot |\phi_{ij}|^2 \right) \left( \sum_{k=1}^p \text{EMAC}_{ik}^I \cdot |\phi_{ik}|^2 \right)}{\sum_{j=1}^q |\phi_{ij}|^2 \cdot \sum_{k=1}^p |\phi_{ik}|^2} \end{aligned} \quad (14)$$

where  $\phi_{ij}$  and  $\phi_{ik}$  are mode shape and modal participation components, respectively. A weighting factor of  $|\phi|^2$  is used to achieve an energy-type emphasis.

In the original formulation of ERA,<sup>5</sup> a parameter known as the modal amplitude coherence ( $\gamma$ ) was introduced. EMAC is a much improved version of this formulation. Under certain common conditions,  $\gamma$  values can be high (even 100%) for all eigenvalues. This insensitivity is avoided with EMAC. The term "extended" in the new name refers to the extension of the primary data analysis window for the final block row and column of the generalized Hankel matrix. EMAC quantifies

the temporal consistency of the identified modal parameters by measuring the predictability of the results in this extended time interval. This test is much more difficult to fulfill than testing only the predictability of the results in the primary analysis window, which is what  $\gamma$  does. EMAC also involves many fewer calculations than  $\gamma$ . Most importantly, however, it provides a more sensitive, more reliable approach with proven usefulness based on many successful applications.

### Modal Phase Collinearity

MPC quantifies the spatial consistency of the identification results. For classical normal modes, all locations on the structure vibrate exactly in-phase or out-of-phase with one another; i.e., the corresponding mode shape is a real or "monophase" vector.

With monophase behavior, the variance-covariance matrix of the real and imaginary parts of the mode shape vectors has only one nonzero eigenvalue. If the identified mode shape phase angles are uncorrelated, on the other hand, the two eigenvalues of this matrix will be approximately equal. MPC quantifies the degree of monophase behavior by comparing the relative size of the eigenvalues of the variance-covariance matrix.

Let  $\Phi_i'$  and  $\Phi_i''$  be the real and imaginary parts, respectively, of the identified mode shape for mode  $i$ . Calculate the variance and covariance of the real and imaginary parts

$$\begin{aligned} S_{xx} &= \Phi_i'^T \Phi_i' \\ S_{yy} &= \Phi_i''^T \Phi_i'' \\ S_{xy} &= \Phi_i'^T \Phi_i'' \end{aligned} \quad (15)$$

Letting

$$\eta = \frac{S_{yy} - S_{xx}}{2S_{xy}} \quad (16)$$

the eigenvalues of the variance-covariance matrix are

$$\lambda_{1,2} = \frac{S_{xx} + S_{yy}}{2} \pm S_{xy} \sqrt{\eta^2 + 1} \quad (17)$$

MPC for mode  $i$  is then defined as follows:

$$MPC_i = \left( \frac{\lambda_1 - \lambda_2}{\lambda_1 + \lambda_2} \right)^2 \quad (x \ 100\%) \quad (18)$$

$MPC_i$  values range from 0 for a mode with completely uncorrelated phase angles to 100% for a monophase result.

This formulation is based on the original definition of modal phase collinearity ( $\mu$ ),<sup>5</sup> except normalized to generate values ranging from 0–100%. The smallest possible value of  $\mu$  was inadvertently limited to 25%. Also, the following two practical extensions of the MPC concept are presented for the first time.

### Unweighted MPC

The definition of MPC in Eqs. (15–18) provides a natural weighting based on the magnitude of the individual mode shape components. This weighting is desirable because phase angle results for small experimentally determined mode shape components are often inaccurate due to measurement limitations. However, it is also useful to repeat the calculations without this natural weighting imposed; i.e., by normalizing each mode shape component to unit magnitude before calculating the variance and covariance values. For global modes, this "unweighted MPC" will be approximately equal to the standard weighted MPC discussed above. However, for local modes, the unweighted value will be substantially smaller than the weighted value. The magnitude of the difference provides

a quick and effective indicator of global versus local response behavior.

### Phase Rotation for Free-Decay Data

When ERA is applied to displacement/force or acceleration/force impulse response functions, the identified mode shapes display the classical characteristic of large imaginary parts and small real parts. With free-decay data for arbitrary initial conditions, however, the identified mode shapes have arbitrary mean phase angle. Before these shapes can be plotted as geometric deformations or used in certain subsequent calculations such as the phase resonance criterion,<sup>9</sup> they must be rotated to align best with  $\pm 90$  deg. The necessary rotation angle  $\alpha$  is determined as

$$\alpha = \tau + \pi/2 \quad (19)$$

where

$$\tau = \tan^{-1} \left[ \eta + \text{sgn}(S_{xy}) \sqrt{\eta^2 + 1} \right] \quad (20)$$

## Application Example

### Test Article and Modal Survey

Figure 1 shows the controls-structures-interaction (CSI) evolutionary model (CEM)<sup>8</sup> located at the NASA Langley Research Center. The structure is part of a testbed used to develop CSI ground test methods. It has been designed to possess the dynamic properties of a typical future large spacecraft. These properties include low frequency structural modes, modal clusters, local appendage dynamics, and both high and low levels of damping.

The CEM consists of a 55 ft-long truss with several appendages that possess varying degrees of flexibility. The truss is constructed of aluminum tubes assembled into 10-in. cubical bays. The structure is suspended by two cables attached to the laboratory ceiling through isolation springs. By using soft springs in series with long suspension cables, the six "pseudo-rigid body" modes have frequencies below 1 Hz. The first flexible mode occurs at 1.5 Hz with approximately 30 modes occurring below 10 Hz. For modal identification experiments, the CEM was instrumented with 195 piezofilm accelerometers and 18 servo accelerometers. Excitation was supplied through 16 on-board cold-gas thrusters operating in pairs at 8 locations. The thrusters produce up to 2.2 lbs of force over a bandwidth of approximately 45 Hz.

A modal test was performed using all 213 accelerometers and 8 thruster pairs in a multi-input, multi-output (MIMO) test configuration. Uncorrelated shaped random noise signals

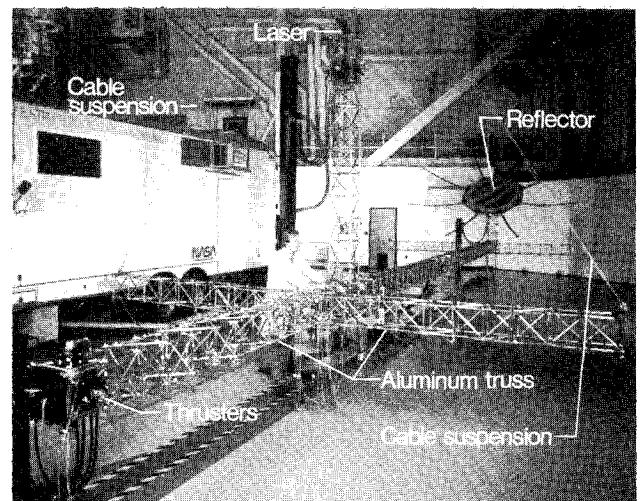


Fig. 1 CSI evolutionary model.

were applied simultaneously to all exciters. A total of 1704 ( $8 \times 213$ ) frequency response functions (FRFs) with 2048 lines of resolution from 0–50 Hz were generated. The FRFs were computed using a commercial MIMO FRF calculation procedure<sup>10</sup> with 100 averages to minimize noise effects. Frequency lines from 0–6.25 Hz were extracted and inverse Fourier transformed to obtain impulse response functions for the ERA analyses discussed later.

Before performing the ERA analysis, it is useful to calculate the average power spectrum (APS) and mode indicator function (MIF)<sup>11</sup> directly from the FRF data. These functions are defined as follows:

$$\text{APS}(f) = \frac{\sum_{i=1}^N |H_i(f)|^2}{N} \quad (21)$$

$$\text{MIF}(f) = \left[ 1.0 - \frac{\sum_{i=1}^N |H_i(f)| |H_i(f)|}{\sum_{i=1}^N |H_i(f)|^2} \right] \cdot 1000 \quad (22)$$

where  $H_i(f)$  and  $|H_i(f)|$  are the real part and magnitude, respectively, of the  $i$ th FRF.  $N$  is the total number of FRFs included in the calculation. Both of these functions display peaks at each natural frequency. Additionally, the APS shows the relative magnitude of each modal response. The MIF provides no information concerning modal response magnitudes; however, the resolution of individual modes is much higher.

The APS and MIF calculated using all 1704 FRFs are shown in Fig. 2. It can be stated with certainty based on these results that at least 16 modes occur between 0 and 6.25 Hz. However, because of the complexity of the structure and indicated by NASTRAN analytical predictions,<sup>8</sup> some of the observed peaks probably represent more than one mode. For example, two pseudorigid body modes are known to occur between 0.7 and 0.8 Hz. Also note the wide variation of modal response magnitudes indicated by APS peaks ranging from 100 to less than 1. Modes with low response magnitudes such as those near 0.2 and 3.4 Hz are much more difficult to identify accurately than modes with high response magnitudes.

#### Identification Results

Previous experience has shown that significant changes can occur in ERA analyses as a function of the assumed number of modes. Optimum accuracy for different modes typically occurs at different numbers of assumed modes. Also, weakly excited modes often require relatively high numbers of assumed modes to be properly identified. For these reasons, the assumed number of modes is incremented over a wide range

of values in most applications. Based on the estimated number of modes from Fig. 2, the assumed number of modes is incremented for the CEM data analysis from 2 up to 60 (in steps of 2).

The identified natural frequencies as a function of the assumed number of modes are plotted in Fig. 3. These results were generated using all 1704 FRFs simultaneously in an  $8 \times 213$  MIMO analysis. Each row of results corresponds to a separate ERA analysis with a specified number of assumed modes. Each detected mode is represented by a vertical dash at the associated frequency. The confidence in each result is expressed by the length of the vertical dash, which is proportional to CMI. The highest confidence (CMI = 100%) is attained if the distance between minor tic marks on the vertical axis is filled. Eigenvalues with low CMI (less than 5%) are excluded from the figure.

A wide spread of CMI values is indicated by the presence of both solid and dashed or dotted lines. Several modes with low CMI occur in the frequency interval from 2.6–3.3 Hz. These modes are associated with the flexible reflector and are known to occur in clusters. No thrusters are located directly on the reflector surface so that these modes cannot in general be excited individually. The result is a reduction of controllability and a corresponding decrease in both temporal and spatial consistency. Although engineering judgment is required to determine the cause of low CMI values, low values reliably indicate those results that should not be accepted verbatim as accurate modal parameters. As illustrated later in this paper, results for particular modes can often be improved without performing additional tests once this initial CMI information is available. Of course, additional tests can also be performed once certain areas of difficulty are known.

Figure 4 shows an expanded view of the results in a narrow frequency interval near 0.9 Hz. Corresponding damping, EMAC, and MPC results are also shown. This example illustrates a typical degree of accuracy variation that occurs as a function of the assumed number of modes. Although damping estimates range from 0.5–2.5%, the highest confidence is associated with values between 0.5–1.0% because of the corresponding high EMAC values. MPC values for this mode show only small variations indicating stable mode shape identification.

A more complicated situation is shown in Fig. 5 for the mode (or modes) near 4.1 Hz. The highest EMAC and MPC values occur near 30 assumed modes. At higher numbers of assumed modes two additional eigenvalues are identified. These additional results may be weakly excited and/or weakly observed local modes such as the vibration of the rubber air-supply hoses for the thrusters. The important thing to note is that their presence causes significant perturbations in both EMAC and MPC for what is believed to be a single structural

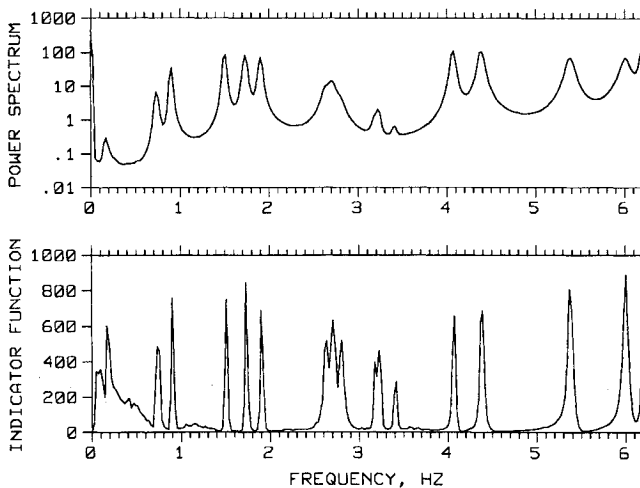


Fig. 2 Average power spectrum and mode indicator function.

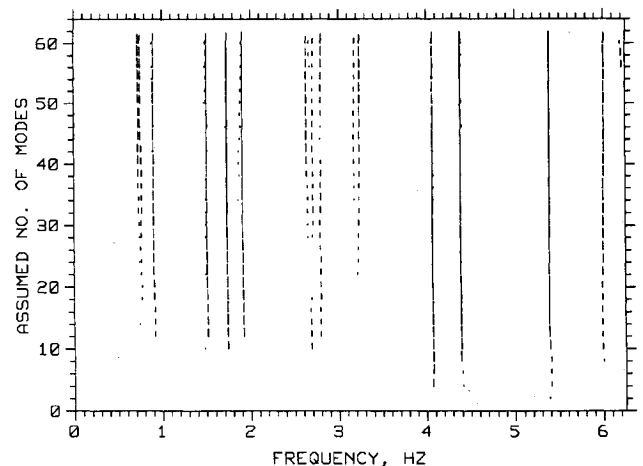


Fig. 3 Identified natural frequencies (lengths of vertical dashes proportional to CMI).

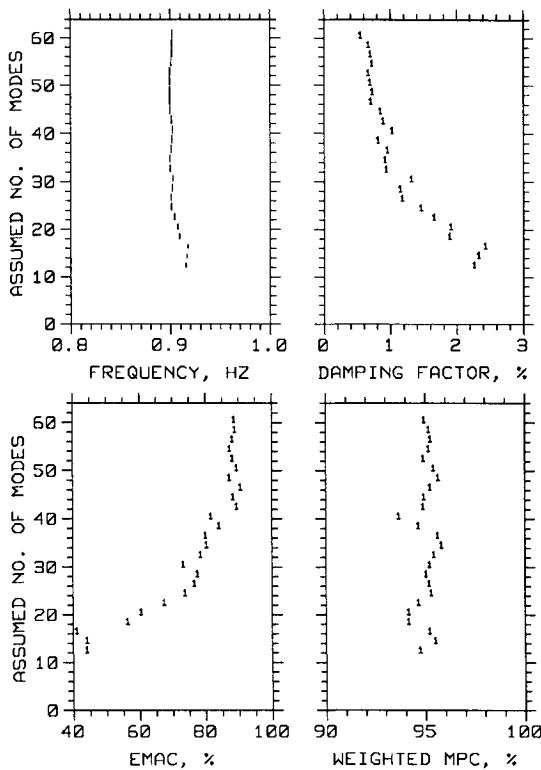


Fig. 4 Identification results at 0.9 Hz.

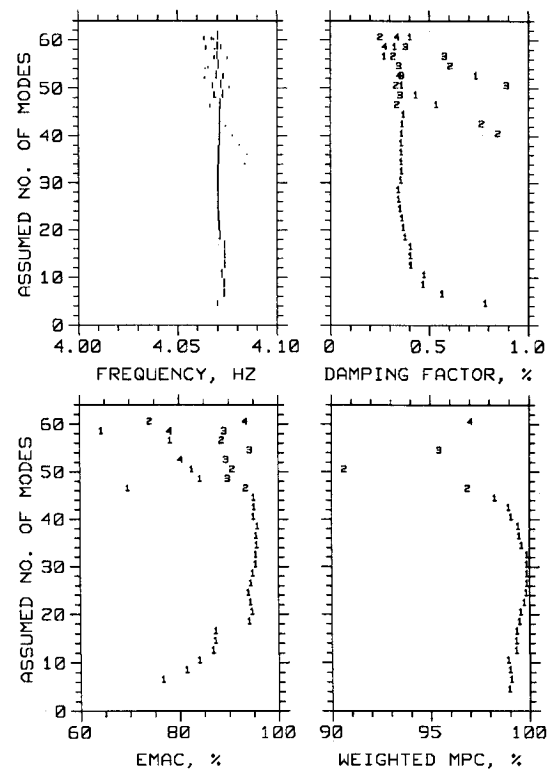


Fig. 5 Identification results at 4.1 Hz.

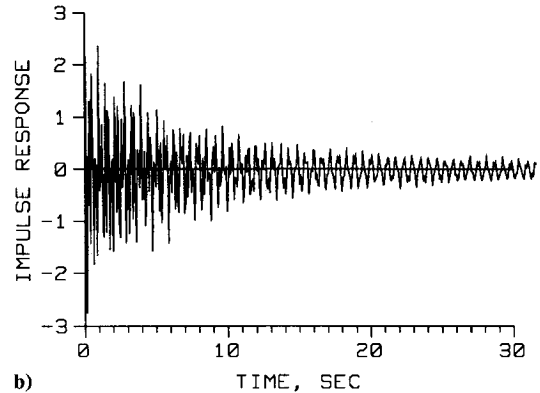
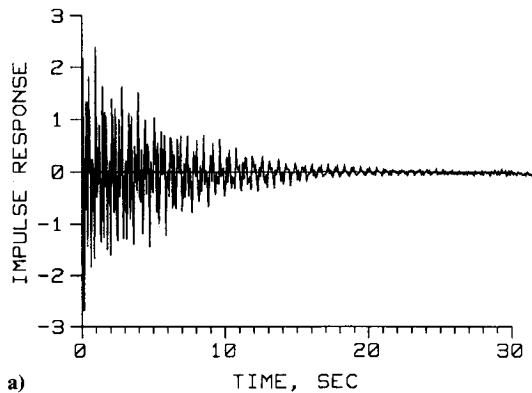


Fig. 6 Comparison of measured and reconstructed responses: a) measured and b) reconstructed.

mode in this frequency interval. Again, the EMAC and MPC results (and their product, CMI) reliably indicate those results with highest confidence based on both temporal and spatial consistency calculations.

Figure 6 compares measured and reconstructed impulse response functions. The ERA analysis used the first 4.5 s of measured data. There is only a slight difference between the two functions in this interval. Beyond the 4.5 s data analysis interval, however, the differences between the two functions become considerably larger. Although it is difficult to know the exact reason for the discrepancy between measured and reconstructed responses beyond the data analysis interval, EMAC values in this situation would decrease, highlighting the inconsistency. EMAC quantifies the temporal consistency of each mode beyond the primary data analysis interval on an input-by-input and measurement-by-measurement basis.

Based on CMI (the product of EMAC and MPC), the best results for the 4.1 Hz mode (Fig. 5) occurred using 30 assumed modes. The corresponding mode shape is shown in Fig. 7a. This result is highly believable based on the smoothness and uniformity of motion over the large set of 213 response measurements. Furthermore, this shape and the corresponding fre-

quency are well predicted by a NASTRAN finite element analysis.<sup>8</sup> Although CMI for this mode is high indicating a reliable result based on consistency calculations, it is also important to examine the physical deformation pattern of the mode to achieve complete confidence.

To illustrate the rationale for using the standard weighted MPC in the CMI calculation rather than the alternative unweighted MPC described earlier, Figure 7b shows the distribution of magnitude and phase results for this 4.1 Hz mode. The values are arranged in descending order of magnitude. Overall, the identified phase angles cluster near the ideal values of  $\pm 90$  deg, and MPC is accordingly high. However, a trend of increasing phase angle scatter from  $\pm 90$  deg at smaller magnitudes is seen. These relatively large phase angle errors at small magnitudes are attributed to finite measurement resolution. The overall mode shape is clearly accurate, however, including accurate identification of small magnitudes at appropriate locations (based on Fig. 7a). If unweighted MPC values are used instead of weighted values in the CMI computation, the CMI value would be unnecessarily lowered.

It is important to realize that MPC decreases from the ideal value of 100% for many reasons other than low response

magnitude. For example, significant phase angle errors can occur for closely spaced modes that are inadequately uncoupled by excitation. For structures, with large numbers of closely spaced modes, such as mini-mast, which had 108 modes between 14 and 22 Hz due to the bending of individual truss members,<sup>7</sup> it may be impractical to apply a sufficient number of exciters. Another cause of lower MPC values is the occurrence of true "complex modes" whose eigenvectors are, in fact, significantly nonmonophasic (i.e., complex) due to a nonproportional damping distribution. Although this situation can occur in practice, low MPC values are more often the result of identification difficulties, many of which can be eliminated once they are detected and understood.

#### Improvement of Results

With complex structures such as CEM, many modes involve significant motion at only a fraction of the total number of measurement locations. In such cases, identification accuracy can be improved by de-emphasizing data with low consistency. Both input and output EMAC values are examined to determine optimum excitation and measurement locations for particular modes based on the results of an initial ERA analysis.

In this application, weighted EMAC values (i.e., EMAC multiplied by the square of the corresponding mode shape coefficient) are used to select optimum excitation and measurement locations. The analyses use data only for the selected exciters. All response measurements are included in the analysis, however, so that complete mode shape information is obtained. Additional emphasis is achieved for the target mode by retaining only the selected response locations in all block rows of the generalized Hankel matrix below row  $q$ , where  $q$  is the total number of response measurements. The theoretical basis for permitting the deletion of rows (or columns) of data in the

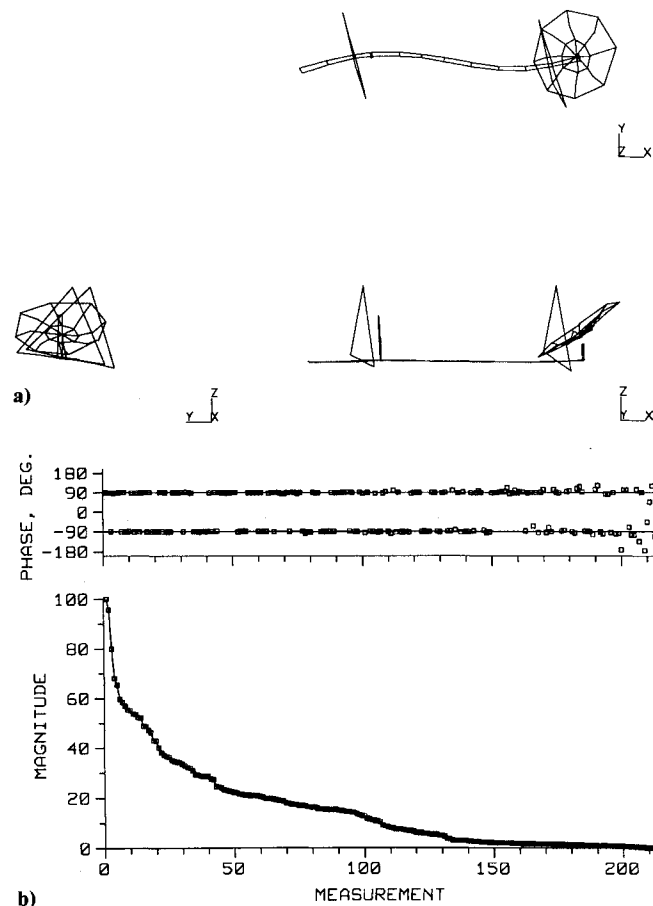


Fig. 7 Identified mode at 4.1 Hz using 30 assumed modes: a) mode shape and b) distribution of magnitude and phase values.

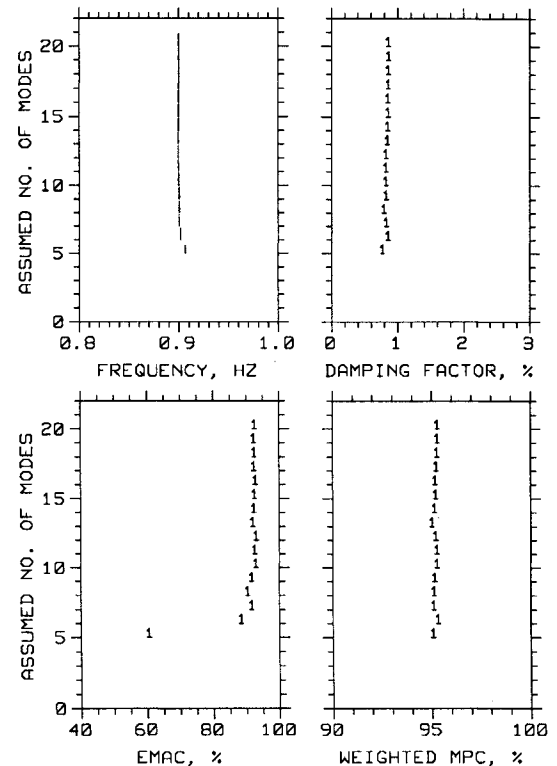


Fig. 8 Improved identification of 0.9 Hz mode.

ERA Hankel matrices without affecting the rank of the matrix is well established.<sup>5,12</sup> The approach discussed here is the authors' practical implementation of the theory.

An example of the improvement achieved using this approach is shown in Fig. 8. These results for the 0.9 Hz mode extend only up to 20 rather than to 60 assumed modes as before (a higher number of assumed modes was unnecessary). Comparing Fig. 8 with Fig. 4 over the range of 2–20 assumed modes, significant improvement is observed. EMAC values are appreciably higher and more stable. Corresponding frequency and damping results also show improved stability. Although MPC fluctuates much less than in Fig. 4, the average value of 95% is still obtained. Lack of improvement of MPC was unexpected. Normally, such stability is not obtained at MPC values less than 99%. This behavior is attributed to the use of piezofilm accelerometers in the experiment. These sensors have considerable phase shift at 0.9 Hz, which is near the lower limit of their operating range. Although all 195 piezofilm sensors were of the same model, appreciable differences in phase response occurred at 0.9 Hz among the individual units. This anomaly did not occur at higher frequencies.

#### Conclusions

The concepts discussed in this paper have been developed over the course of several years in conjunction with many applications of the eigensystem realization algorithm. The result of these efforts is a single parameter known as the consistent-mode indicator, which reliably indicates the relative confidence of each identified mode on the basis of both temporal and spatial consistency calculations. In practice, modes with indicator values greater than approximately 80% can generally be accepted verbatim. Modes with lower values, however, require additional attention. By examining both the extended modal amplitude coherence, which measures temporal consistency, and the modal phase collinearity, which measures spatial consistency, an explanation for low indicator values can often be developed. Once an explanation is obtained, additional tests and/or data analyses can be performed to improve these results. One approach for improving results of selected

modes by de-emphasizing data with low consistency was illustrated.

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# Fundamentals of Tactical and Strategic Missile Guidance

Paul Zarchan  
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